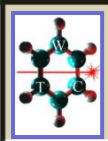


Optical fiber interferometry

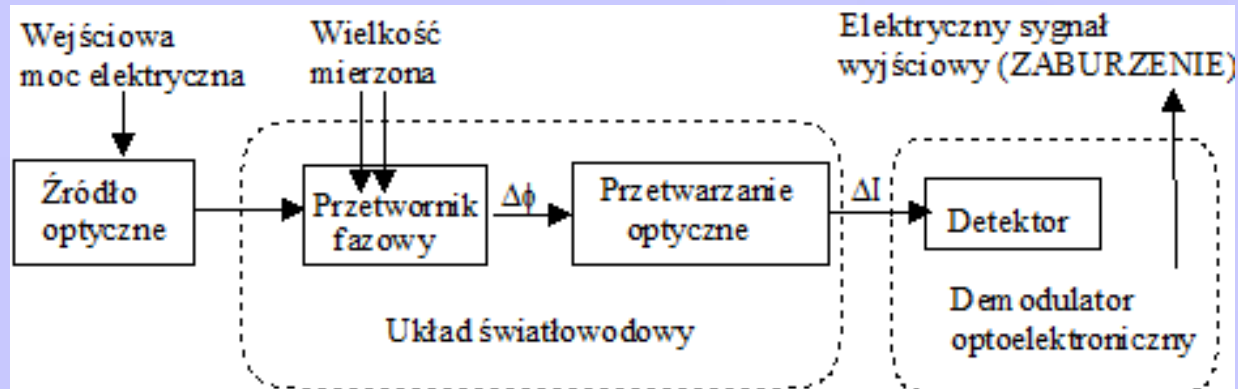
- **Introduction**
- **General principles**



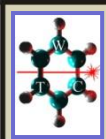


Introduction

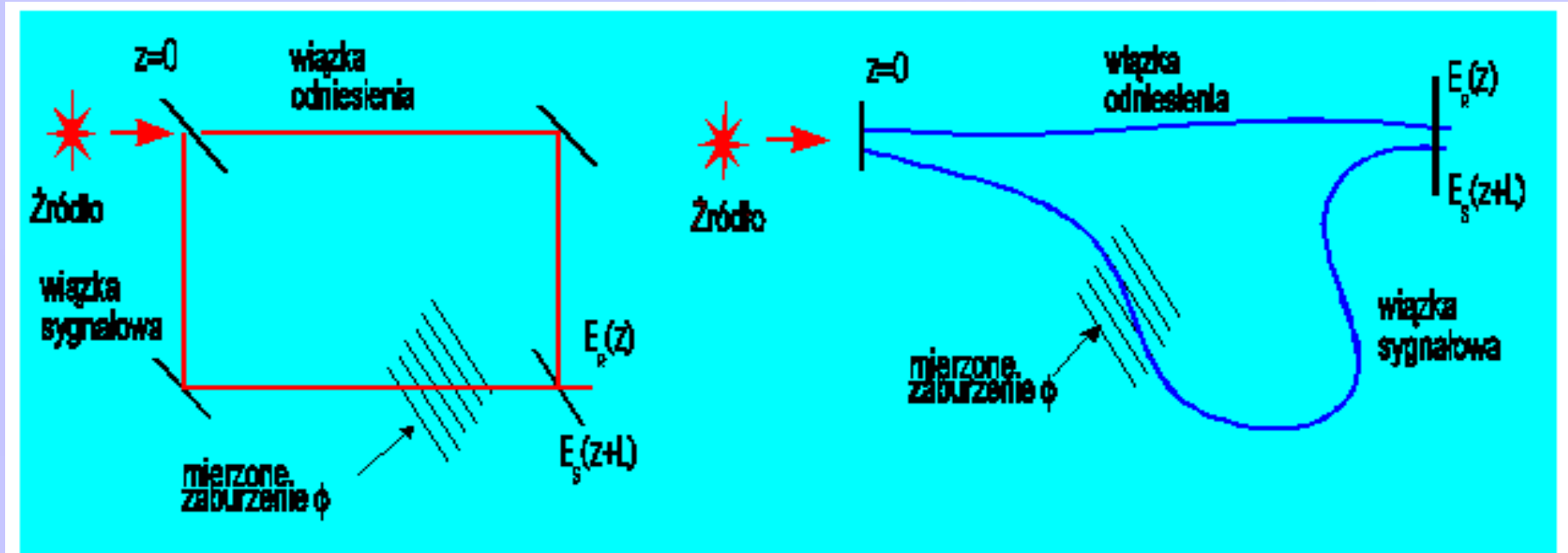
They are a new sensor class, where general principle of operation based on changes by measurand the optical way or polarization properties of optical fiber. Such sensors are named *intrinsic* or *phase* sensors.

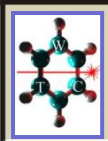


For optical fiber, optical processing technique which process phase information on intensity information (ΔI), is called *optical-fiber interferometer*. It is a simple fiber-optic device allows observation of a interference between two or more optical beams.



Fiber optic interferometers are equivalents all well-known optical bulk interferometers, where light is closed in structure of single-mode or polarization preserving optical fiber.



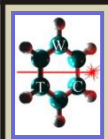


General principles

An optical sensor may be formally defined as a device in which an optical signal is modulated in response to a measurand field. Let us assume that the source has some well-defined wavelength spectrum, and that the electric field at wavelength λ is $\mathbf{E}(\lambda)$, in unit bandwidth. Then the corresponding received electric field will be:

$$\mathbf{E}'(\lambda) = \mathbf{T}(\mathbf{X}, \lambda)\mathbf{E}(\lambda)$$

where $\mathbf{T}(\mathbf{X}, \lambda)$ is the propagation matrix describing the sensing element and \mathbf{X} is a vector describing the physical environment, including terms representing temperature, stress, electromagnetics fields. The function of the signal processing used in the sensor system is to invert above relation, to find \mathbf{T} , invert again to find \mathbf{X} , and then to identify and evaluate the relevant component(s) of \mathbf{X} to recover the desired measurand.



It is instructive to express T as a product of terms, each describing a physically observable effect on the transmitted beam, such that:

$$T = ae^{i\phi_1} B$$

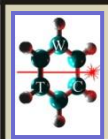
where: a – scalar transmittance, ϕ_1 – the mean phase retardance, and B are all both dispersive and environmentally sensitive. In a multimode fiber, spatial coherence is not fully maintained, so that for many simple multimode fiber sensors the device is entirely characterized by the dispersive scalar transmittance, a .

In monomode systems, sensing mechanisms based on modulation of any one or combination of the parameters a , ϕ_1 , B . In practice, the transmittance of SMF shows only weak environmental sensitivity, and intrinsic monomode sensors are thus generally based on phase and polarization modulation.

The transfer function of a monomode sensing element in the Jones calculus is:

$$E' = a_0 E e^{i\phi_1} B$$

where: a_0 – scalar constant (transmittance) and B – the Jones matrix.



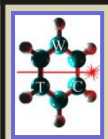
For a SMF possessing perfect cylindrical symmetry $\mathbf{B} = \mathbf{I}$, but general it is necessary to consider the effects of birefringence within the fiber. For example for a linear birefringent fiber:

$$\mathbf{B} = \mathbf{B}_l = \begin{bmatrix} e^{i\phi_2/2} & 0 \\ 0 & e^{-i\phi_2/2} \end{bmatrix}$$

Such a fiber is characterised by two linear polarization eigenmodes, such that ϕ_2 is the induced relative phase retardance between the eigenmodes caused by propagation through the fiber. For a circularly birefringent fiber:

$$\mathbf{B} = \mathbf{B}_c = \begin{bmatrix} \cos\phi_3 & -\sin\phi_3 \\ \sin\phi_3 & \cos\phi_3 \end{bmatrix}$$

where: $2\phi_3$ – the induced relative phase retardance between the eigenmodes, which in this case are left and right circularly polarized states.



The environmental sensitivity of the fiber can be described in terms of above phase ϕ_i ($i=1,2,3$) dependencies on external stimuli such as temperature (T), pressure (P), and strain (Δ/l):

$$\frac{\partial \phi_i}{\partial X} = \frac{2\pi}{\lambda} \left(n_i \frac{\partial l}{\partial X} + l \frac{\partial n_i}{\partial X} \right); X = T, P, \Delta l, \dots; i = 1, 2, 3$$

l – the length of the fiber, n_i - relative refractive index of fiber (1- fiber, 2- difference between refractive index of two linear eigenmodes, 3 – difference of refractive index between circular eigenmodes)

Table 10.1 Temperature and strain sensitivity coefficients measured using 0.1 m of highly birefringent monomode optical fiber (York Technology “bow tie” fiber, 3 mm beat length) at a wavelength of 633 nm

X	$(1/l)(\partial \phi_1 / \partial X)$	$(1/l)(\partial \phi_2 / \partial X)$
T	100	5 rad K ⁻¹ m ⁻¹
Δ/l	6.5×10^6	6.5×10^4 rad m ⁻¹

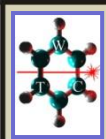
Transfer function:

$$I = I' \{1 + V(\gamma) \cos[\phi - \alpha(\gamma)]\}$$

bulk system

$$I = I' \{1 + V(\gamma, \text{SOP}) \cos[\phi - \alpha(\gamma, \text{SOP})]\}$$

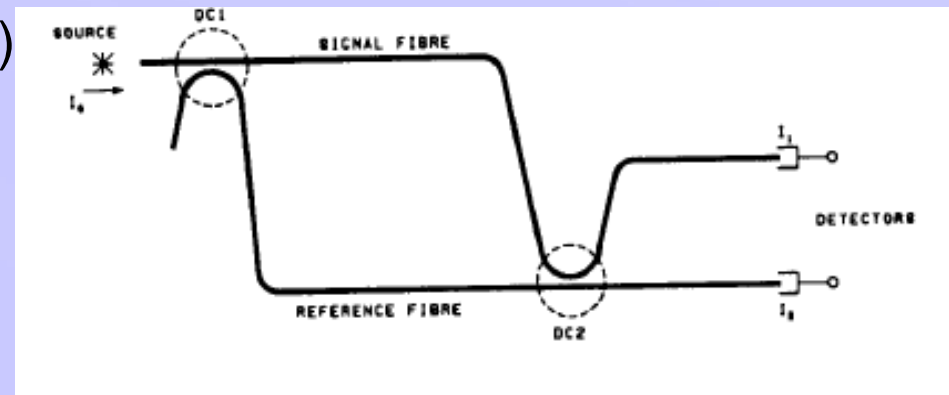
fiber-optic system



Ideal construction (without polarisation)

$$E_1 = k_{2c} \exp(i\phi_b) k_{1c} E_0(\tau_b) + k_{2t} \exp(i\phi_a) k_{1t} E_0(\tau_a)$$

$$E_2 = k_{2t} \exp(i\phi_b) k_{1c} E_0(\tau_b) + k_{2c} \exp(i\phi_a) k_{1t} E_0(\tau_a)$$



Coupler description: $k_{1c} = ik'_{1c}$

Interference: $I_1 = \langle E_1 \cdot E_1^* \rangle$

$$I_1 = k'^2_{1c} k'^2_{2c} \langle E^2_0(\tau_b) \rangle + k^2_{1t} k^2_{2t} \langle E^2_0(\tau_a) \rangle + 2\text{Re}[-k'_{1c} k'_{2c} k_{1t} k_{2t} e^{i(\phi_a - \phi_b)} \langle E_0(\tau_a) \cdot E^*_0(\tau_b) \rangle]$$

The degree of coherence of the source: $\gamma(\tau_a - \tau_b) = \langle E_0(\tau_a) \cdot E^*_0(\tau_b) \rangle / I_0$

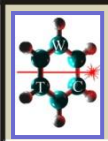
Visibility (contrast) of the interference: $V = (I_{\max} - I_{\min}) / (I_{\max} + I_{\min})$

Intensity on detector: $I_1 = I_0 [1 - V \cos(\phi_a - \phi_b)]$ $I_2 = I_0 [1 + V \cos(\phi_a - \phi_b)]$

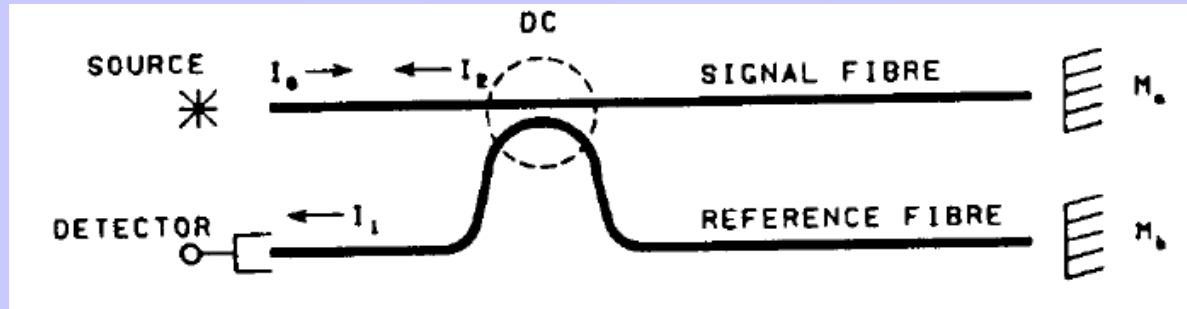
$$V = \frac{2k_{1c} k_{2c} k_{1t} k_{2t}}{k^2_{1c} k^2_{2c} + k^2_{1t} k^2_{2t}} \gamma$$

$$|\gamma(\tau)| = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \exp i [\phi(t + \tau) - \phi(t)] dt$$

Practially: $\tau_c \gg \tau_a - \tau_b$ and $k'_{1c} = k_{1t} = 1/\sqrt{2}$, thus $\gamma=1, V=1$



Ideal construction (without polarisation)



Electrical field on detector:

$$E_1 = k_c m_a \exp(i\phi_a) k_t E_0(\tau_a) + k_t m_b \exp(i\phi_b) k_c E_0(\tau_b)$$

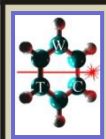
$$E_2 = k_t m_a \exp(i\phi_a) k_t E_0(\tau_a) + k_c m_b \exp(i\phi_b) k_c E_0(\tau_b)$$

Intensity on detector:

$$I_1 = I_0 [1 - V \cos(\phi_a - \phi_b)]$$

$$I_2 = I_0 [1 + V \cos(\phi_a - \phi_b)]$$

$$V = 2m_a m_b \gamma / (m_a + m_b)$$



Dual wavelength interferometry (λ_1 i λ_2) – increase of the dynamic range

Transfer function:

$$I = I_0 \left(1 + \cos \frac{2\pi nl}{\lambda_1} \right) + I_0 \left(1 + \cos \frac{2\pi nl}{\lambda_2} \right)$$

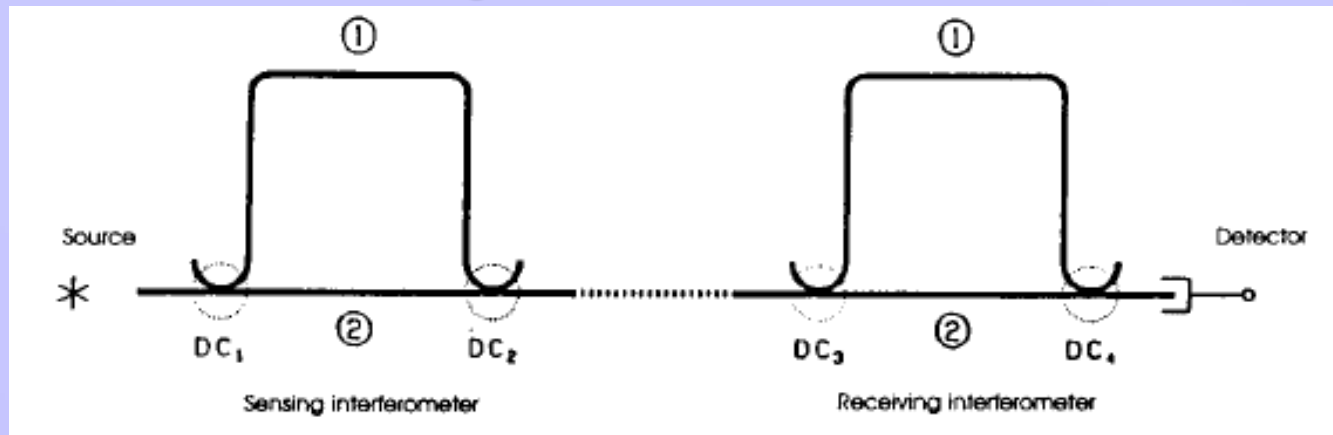
$$I = I_0 \left[1 + V \cos \left(\pi nl \frac{\lambda_1 + \lambda_2}{\lambda_1 \lambda_2} \right) \right]$$

$$V = \cos \left[\pi nl \left(\frac{\lambda_2 - \lambda_1}{\lambda_1 \lambda_2} \right) \right]$$

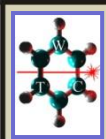
Increase of the dynamic range of the factor

$$\lambda_2 / (\lambda_2 - \lambda_1)$$

White light interferometer



An alternative technique for the extension of the dynamic range concerns the use of short coherence length source with gives possibility to measure phase as well as visibility.



Input electrical field:

$$\mathbf{E} = \mathbf{E}_{11} + \mathbf{E}_{12} + \mathbf{E}_{21} + \mathbf{E}_{22}$$

$$\mathbf{E}_{ij} = \frac{1}{4} E_0(\tau_{jk}) \exp i\omega(\tau_{jk})$$

$$\tau_{11} - \tau_{12} = \tau_{21} - \tau_{22} = \tau_a$$

$$\tau_{12} - \tau_{22} = \tau_{21} - \tau_{11} = \tau_b$$

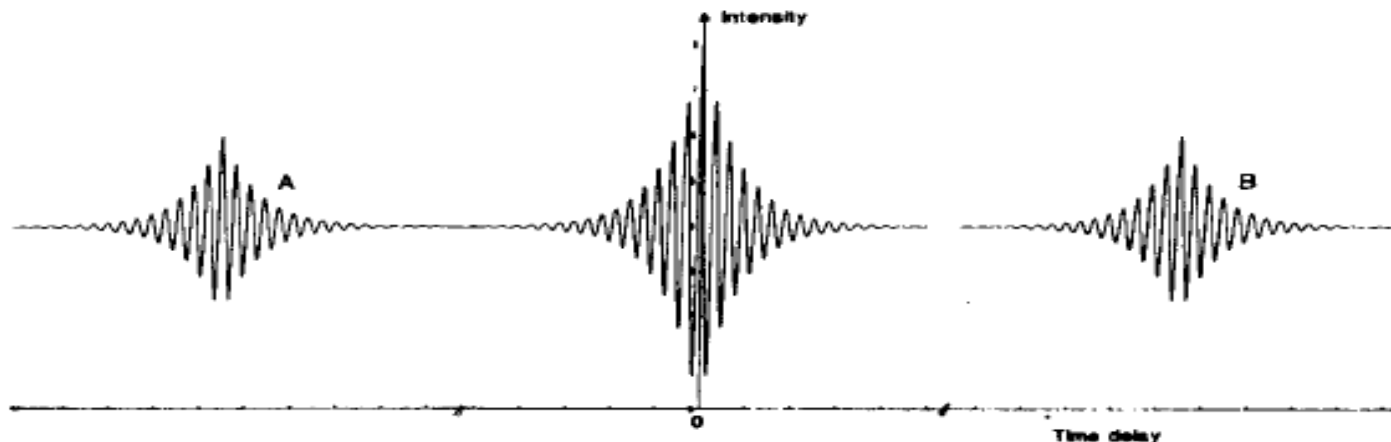
Describe time difference of light propagation between ways 1 and 2 for sensing and receiving interferometers => transfer function:

$$\tau_b \approx \tau_c$$

$$I = \frac{1}{4} I_0 \left[1 + \cos \omega \tau_a \right]$$

$$\tau_a \approx \tau_b$$

$$I \approx \frac{1}{4} I_0 \left[1 + \frac{1}{2} \cos \omega (\tau_a - \tau_b) \right]$$



Transfer Function of Tandem Interferometers

Fig. 10.6b Interference fringes produced by the arrangement shown Fig. 6a, as a function of path length imbalance (and hence time delay) in the receiving interferometer, where the path imbalance of the sensing interferometer is much greater than the coherence length of the source. The group of interference fringes around the origin correspond to near-zero imbalance of the receiving interferometer; groups A and B correspond to the receiving interferometer balancing the path difference in the sensing interferometer, and satisfy the conditions $\tau_a + \tau_b \approx \tau_c$ and $\tau_a - \tau_b \approx \tau_c$ respectively, where τ_a , τ_b and τ_c are as defined in the text.